

宇田方程式	$\int_{-\infty}^{\infty} dt \left\{ \frac{1}{2m} \left[-\frac{i\hbar}{\alpha} \frac{\delta}{\delta \chi(t)} - m \frac{d\chi(t)}{dt} \right]^2 - \frac{m}{2} \left[\frac{d\chi(t)}{dt} \right]^2 \right\} \Phi[\chi] = 0$
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ハミルトニアン	$H = H_0 + H_1 \begin{cases} H_0 \equiv \frac{1}{2m\epsilon\alpha^2} \sum_{k=-\infty}^{\infty} [(P^k)^2 + m\alpha(P^k\chi^k + \chi^k P^k)] \\ H_1 \equiv -\frac{1}{\epsilon\alpha} \sum_{k=-\infty}^{\infty} P^k \chi^{k+1} \end{cases}$
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自由粒子1個の定常状態	$\frac{\left[\left(-i\hbar \frac{d}{dx} \right)^2 - i\hbar m\alpha \left(1 + 2x \frac{d}{dx} \right) \right] u(x) = e u(x)}{(P)^2 + m\alpha(PX + XP)}$
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$$u(x) = \sum_{n=0}^{\infty} a_n x^n \begin{cases} 2\hbar^2 a_2 + (e + i\hbar m\alpha) a_0 = 0 \\ \hbar^2(n+2)(n+1)a_{n+2} + [e + i\hbar m\alpha(2n+1)]a_n = 0 \quad (n \geq 1) \end{cases}$$

$$a_{2n}(e) = \frac{1}{(2n)!} \left(\frac{-4im\alpha}{\hbar} \right)^n \frac{\Gamma(n + 1/4 + e/(4i\hbar m\alpha))}{\Gamma(5/4 + e/(4i\hbar m\alpha))} \frac{e + i\hbar m\alpha}{4i\hbar m\alpha} a_0(e)$$

$$a_{2n+1}(e) = \frac{1}{(2n+1)!} \left(\frac{-4im\alpha}{\hbar} \right)^n \frac{\Gamma(n + 3/4 + e/(4i\hbar m\alpha))}{\Gamma(3/4 + e/(4i\hbar m\alpha))} a_1(e) \quad \Gamma(n+1) = n\Gamma(n)$$